Optical crystallographic properties of the 4-phenyl- and 4-cyclohexylcyclohexanols. By Walter D. Keller and Herbert E. Ungnade. Geology and Chemistry Departments, University of Missouri, Columbia, Mo., U.S.A.
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The cis and trans isomers of 4-phenylcyclohexanol have been shown by thermal analysis to form a system with a single eutectic point in which solid solutions are absent below $10 \%$ and above $90 \%$. The cis and trans isomers of 4-cyclohexylcyclohexanol form a system with two eutectic points which is indicative of compound formation and miscibility (Ungnade, 1948). The unexpected differences in the two pairs of isomers are in agreement with the optical crystallographic properties of the substances.
trans-4-Phenylcyclohexanol (I) is biaxial, positive in character, with 2 V about $75-80^{\circ}$, whereas cis-4-phenylcyclohexanol (II) is uniaxial and negative in optical character. The two isomers necessarily crystallize individually in different crystal systems and result in the formation of a eutectic when mixed. (I) is probably monoclinic and (II) is probably hexagonal, but the crystal morphology is too imperfectly developed to permit definite determination.
cis-4-Cyclohexylcyclohexanol (IV) is biaxial, positive
in character with $2 V$ about $70-75^{\circ}$ and birefringence about 0.038 , whereas trans-4-cyclohexylcyclohexanol (III) is biaxial, with $2 V$ so close to $90^{\circ}$ that the optical character is indeterminate or probably ranges from positive to negative across $90^{\circ}$, and has a birefringence of about 0.026 . Both isomers crystallize in needles. Crystals of (III) are elongated parallel to $X$, those of (IV) are elongated parallel to $Y$. Both crystallize in the monoclinic system and exhibit a degree of similarity in optical character.

The optical properties given in Table 1 were measured with a petrographic microscope. Refractive indices were determined by comparison in white light with standardized immersion aqueous solutions of potassium mercuric iodide. Optic angles were estimated from interference figures well centered in orientation.

## Reference

Ungnade, H. E. (1948). J. Organ. Chem. 13, 361.

Table 1. Optical properties

| Compound | Refractive indices* |  |  | Optical character | $2 V\left({ }^{\circ}\right) \dagger$ | Axial dispersion | Elongation |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $X$ | $\underbrace{}_{Y}$ | $Z$ |  |  |  |  |
| I | 1.549 | 1.595 | 1.688 | + | 75-80 | $r>v$ | - |
| II | $E=1.550$ |  | $O=1.581$ |  | 0 |  | $\stackrel{-}{\square}$ |
| III | 1.542 | 1.555 | 1.568 | -? +? $\ddagger$ | 87-90 | $\mathrm{r}>\mathrm{v}$ § | - |
| IV | 1.534\|| | 1.549 | 1.572 | + | 70-75 | $r>v$ | $\pm$ |

* Precision 0.002 except where noted.
$\ddagger$ The optical character was not distinguishable, for the optic axis interference figures appeared as practically straight bars. Some figures suggested ( + ), others ( - ); it is probably ( - ).
§ If the optical sign is negative.
|| Accuracy 0.004.

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Relative phases of diffraction maxima by multiple reflection. By William N. Lipscomb. School of Chemistry, University of Minnesota, Minneapolis 14, Minnesota, U.S.A.
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The occurrence of 'forbidden' maxima in the diffraction of X-rays (Renninger, $1937 a, b$; Davisson \& Haworth, 1944; Pabst, 1939; Collin \& Lipscomb, 1949) and electrons (Raether, 1932; Darbyshire \& Cooper, 1935) has been observed by several investigators. Most often this phenomenon is due to a second diffraction ( $h_{2} k_{2} l_{2}$ ) following a high-intensity first diffraction ( $h_{1} k_{1} l_{1}$ ). The second diffraction emerges from the crystal in the direction of another diffraction $\left(h_{3} k_{3} l_{3}\right)$ whose indices are related to the first two by the equation $\mathbf{h}_{3}-\mathbf{h}_{1}=\mathbf{h}_{2}$, where the $\mathbf{h}_{i}$ are reciprocal-lattice vectors.* It is the purpose of this note to discuss some of the more obvious factors in connection with one experimental arrangement using X-rays to investigate the determination of the relative phases of ( $h_{2} k_{2} l_{2}$ ) and ( $h_{3} k_{3} l_{3}$ ) when the latter is of relatively low but not zero intensity.

* A more complete statement would include the contribution of the intensity of $\left(h_{3} k_{3} l_{3}\right)$ to $\left(h_{1} k_{1} l_{1}\right)$ because of the relation $\mathbf{h}_{3}-\mathbf{h}_{1}=-\mathbf{h}_{2}$, but this is not important for the present discussion.

Such an arrangement would rotate the crystal about $\mathbf{h}_{1}$ as an axis, with the angle between incident beam and axis of rotation such as to maintain the reciprocal-lattice point ( $h_{1} k_{1} l_{1}$ ) on the sphere of reflection for all angular positions of the crystal. Simultaneous to diffraction from any other reflecting reciprocal-lattice point $\left(h_{3} k_{3} l_{3}\right)$ there would then be diffraction in the same direction from a corresponding reciprocal-lattice point ( $h_{2} k_{2} l_{2}$ ) by the mechanism described. Such simultaneous diffractions would show interference effects depending, at least in part, on how the phase change of ( $h_{3} k_{3} l_{3}$ ) compares with the sum of the phase changes of $\left(h_{1} k_{1} l_{1}\right)$ and $\left(h_{2} k_{2} l_{2}\right)$. The intensities of the diffraction maxima so obtained could be compared with those obtained after slight displacement of the axis of rotation so that ( $h_{1} k_{1} l_{1}$ ), and subsequently ( $h_{2} k_{2} l_{2}$ ), were no longer diffracting. If the phase change of ( $h_{1} k_{1} l_{1}$ ) is considered constant for all orientations of the crystal, the information obtained might be presumed to be related to the relative phases of $\left(h_{2} k_{2} l_{2}\right)$ and $\left(h_{3} k_{3} l_{3}\right)$.

Although experimental results have so far been in-

